Monomeric friction coefficient of 1,2-polybutadienes

Weilian Liu, Yuming Yang and Tianbai He*

Changchun Institute of Applied Chemistry, Academia Sinica, Changchun, People's Republic of China (Received 1 March 1988; accepted 31 March 1988)

The monomeric friction coefficient has been determined for 1,2-polybutadienes from dynamic measurements in the transition zone. These values are, respectively, 7.59×10^{-5} , 3.09×10^{-6} , 5.25×10^{-7} and 2.95×10^{-9} g s⁻¹ for 1,2-polybutadienes with vinyl contents of 91%, 76%, 53% and 14%. The structure-dependent behaviour of the monomeric friction coefficient has been discussed. 1,2-Polybutadienes with higher T_g , i.e. with higher vinyl content, have a larger monomeric friction coefficient. The fact that the monomeric friction coefficient increases with increase in cross-sectional area of the polymeric chains describes the effect of chain flexibility on the monomeric friction coefficient. The quantitative relationship between the shortest relaxation time, τ_A , and the tube diameter, d, in Doi-Edwards theory has also been reported: log $\tau_A = -17.24 + 0.38d$.

(Keywords: 1,2-polybutadiene; monomeric friction coefficient; Doi-Edwards theory; cross-sectional area of polymer chains; relaxation time)

INTRODUCTION

One acceptable molecular theory of viscoelasticity for entangled linear polymer liquids is Doi–Edwards theory¹, which is well known as the tube or reptation model². Experimentally, the viscoelastic properties of high-molecular-weight polymers and their concentrated solutions exhibit more complex behaviour than the tube and reptation models predict. Thus, some modifications, e.g. contour-length fluctuation^{3,4} and tube renewal^{4,5}, attempt to bring the theoretical model into better agreement with experiment. However, the basic parameters in Doi–Edwards theory are still used in those modifications. Our interest lies in being able to correlate the structural parameters of polymers with the basic parameters in Doi–Edwards theory and its modifications.

There are two basic parameters in Doi-Edwards theory. One is the monomeric friction coefficient, which characterizes the local dynamics, and the other is the step length of the primitive path, or the tube diameter, which characterizes the topology of the medium. Doi and Edwards take the primitive path step length and the tube diameter to be equal¹, although the latter is thought to be irrelevant except for discussion purposes⁴. Some methods have been given to evaluate the tube diameter from structural parameters of polymers. Graessley⁶ expresses the tube diameter as:

$$d^{2} = \frac{4}{5} \left(\frac{r^{2}}{M} \right) \left(\frac{cRT}{G_{N}^{0}} \right)$$
(1)

where c is the polymer concentration (wt/mole), R the universal gas constant, G_N^0 the plateau modulus, T the temperature and r^2 the mean-square end-to-end distance of molecular weight M. The entanglement spacing:

$$R_{\rm c} = N_{\rm c}^{1/2} l_0 \tag{2}$$

where l_0 is average backbone length of a repeat unit and N_c the critical number of entangled repeating units, is also

considered to be equivalent to the tube diameter^{7,8}. He and Porter give a simple expression for the tube diameter:

$$d = 5.3A^{1/2} \tag{3}$$

where A is the cross-sectional area of the polymer chains.

The monomeric friction coefficient is usually deduced from dynamic measurements in the transition region of polymers. Few expressions for evaluation of the monomeric friction coefficient from structural parameters of polymers have been reported. In this paper, an attempt is given to set up a relationship between the monomeric friction coefficient, which is determined from dynamic measurements, and structural parameters of polymers, e.g. the cross-sectional area of polymer chains, for 1,2polybutadienes.

EXPERIMENTAL

1,2-Polybutadienes with various vinyl contents are the products of anionic polymerization. Microstructure was determined from the i.r. spectrum. Number-average molecular weight was measured with a Knauer membrane osmometer. Molecular-weight distribution obtained from g.p.c. for all samples is less than 1.3. Glass transition temperatures were determined with a linear dilatometer. The density of the samples at 20°C, measured with a density gradient column, was used to calculate the density at 25°C combining with the volume expansion coefficient above T_g of the samples. The characteristic parameters of the 1,2-polybutadiene samples are listed in *Table 1*.

The storage and loss moduli of the samples, E' and E'', were measured with a Rheovibron Dynamic Viscoelastometer model DDV-II-EA. The moduli at four frequencies were measured from a temperature 30°C lower than T_g to a temperature 50°C higher than T_g at a heating rate of 2°C min⁻¹. The strain amplitude of the measurement is 0.02%. The accuracy of the moduli values was checked by repeating measurements twice at each frequency. Except for the glassy region the reproducibility

^{*} To whom all correspondence should be addressed

Table 1 Characteristic parameters of 1,2-polybutadienes

Microstructure				
(vinyl/cis-1,4/trans-1,4)	91/2/7	76/12/12	53/17/30	14/62/24
$M_{\rm n} \times 10^{-3} ({\rm g mol^{-1}})$	145	160	392	122
$T_{g}(\mathbf{K})$	265	246	220	179
$\rho_{25^{\circ}C} (g \text{ cm}^{-3})$	0.884	0.885	0.888	0.889



Figure 1 The temperature dependence of the storage modulus at four frequencies for 1,2-polybutadiene with 91% vinyl content. Open and closed circles refer to different measurements. Data are reproducible at the transition zone

of the moduli is satisfactory (*Figure 1*). The non-reproducibility of the moduli in the glassy region in *Figure 1* is due to non-constant pre-tension with manually mounted samples.

RESULTS AND DISCUSSION

The shift factors for the master curve of 1,2-polybutadiene can be expressed by the WLF equation. The general method to obtain the constants C_1 and C_2 in the WLF equation for 1,2-polybutadiene with various vinyl contents has been reported¹⁰ as:

$$\log \alpha_T = \frac{5.78(T - T_g - 55)}{94.8 + (T - T_g - 55)} \tag{4}$$

The C_1 and C_2 values for the tensile moduli master curve of 1,2-polybutadiene samples in this work (*Table 2*) are determined with:

$$C_1 = \frac{5.78(C_2 + T_g - 298)}{C_2} \tag{5a}$$

$$C_2 = 94.8 - T_s + 298$$
 (5b)

$$T_{\rm s} = T_{\rm s} + 55 \tag{5c}$$

where 94.8, 5.78 and 55 are the constants in equation (4), 298 the reduced temperature for the master curve and T_g the glass transition temperature of the samples. One of those master curves is shown in *Figure 2*. From the master curves of both storage and loss moduli, the method of Ferry and Williams¹¹ is used to obtain the relaxation spectra for 1,2-polybutadiene samples:

$$H(\tau) = AG' d \log G'/d \log \omega \Big|_{1/\omega = \tau}$$
(6)

$$H(\tau) = BG''(1 - |\operatorname{d}\log G''/\operatorname{d}\log \omega|)|_{1/\omega = \tau}$$
⁽⁷⁾

where G' and G'' are taken as one-third of the values of E' and E'', respectively.

The relaxation spectra for all four 1,2-polybutadiene samples are shown in *Figure 3*. With decreasing vinyl content, the position of the transition shifts to shorter times, traversing nearly five decades of the logarithmic timescale. Similar to the methacrylate series¹², 1,2-polybutadiene with lower T_g , i.e. with lower vinyl content, is associated with shorter relaxation time for the transition zone.



Figure 2 Storage modulus master curve reduced to 298 K for 1,2polybutadiene with 91% vinyl content. The C_1 and C_2 values for the shift factor are obtained with equation (5) in the text



Figure 3 Relaxation spectra for 1,2-polybutadienes. The figures refer to vinyl content. Closed circles obtained from storage modulus and open circles obtained from loss modulus

The monomeric friction coefficients for 1,2polybutadiene samples are calculated with:

$$\log \zeta_0 = 2 \log H + \log \tau + \log(6/kT) + 2 \log(2M_0 \pi/a\rho N_A)$$
(8)

where a is the root-mean-square end-to-end length per square root of the number of monomer units, ρ the density of the polymer and M_0 the monomer molecular weight. The experimental relaxation spectra for all four 1.2-polybutadiene samples are matched to the theoretical slope of -1/2 only by drawing a tangent. In equation (8), $\log H = 5.75$ at $\log \tau = -3.67$, $\log H = 5.81$ at $\log \tau =$ -5.17, log H = 5.76 at log $\tau = -5.83$ and log H = 5.54 Pa at $\log \tau = -7.62$ s are, respectively, used for 1,2polybutadienes with vinyl contents of 91%, 76%, 53% and 14%. From the values of unperturbed end-to-end distance for 1,2-polybutadiene in the literature¹³, the vinyl-content-dependent unperturbed end-to-end distance is:

$$(r_0^2/M)^{1/2} = 1.033 - 0.0598X_{1,2} (\text{Å mol}^{1/2} \text{g}^{-1/2})$$

(9)

with correlation coefficient -0.957. Thus, the values of a in (8) are obtained with:

$$a \equiv (r_0^2/M)^{1/2} = (r_0^2/N)^{1/2} (M_0)^{1/2}$$

The monomeric friction coefficients ζ_0 for 1,2polybutadiene are listed in *Table 2*. The monomeric friction coefficient for 1,2-polybutadiene with vinyl content of 91% in *Table 2*, 7.59 × 10⁻⁵ g s⁻¹, is very close to the reported value of 7.76×10^{-5} g s⁻¹. But the monomeric friction coefficient for polybutadiene with vinyl content of 14% in *Table 2* is two orders of magnitude lower than that for 1,2-polybutadiene with 8% vinyl content¹⁴⁻¹⁶.

The significance of the monomeric friction coefficient is the average force per monomer unit required for a chain segment to push its way through its local surroundings at unit velocity. Thus, the monomeric friction coefficient depends on intermolecular forces and chain flexibility of polymer chains. It has long been recognized¹⁷ that a high glass transition temperature is characteristic of a polymer with strong intermolecular forces and high chain stiffness. Therefore, it is not unexpected that 1,2-polybutadiene with higher T_g , i.e. with higher vinyl content, corresponds to higher monomeric friction coefficient. The logarithmic values of the monomeric friction coefficient (in g s⁻¹) and T_g (in K) satisfy:

$$\log \zeta_0 = -17.29 + 0.049T_g \tag{10}$$

with correlation coefficient -0.992.

In the case of the 1,2-polybutadiene series, chain flexibility is the sole factor that determines the glass transition temperature¹³. The nature of the relationship in equation (10) indicates the dependence of the

monomeric friction coefficient on the chain flexibility of 1,2-polybutadienes. One structural parameter describing chain flexibility is the cross-sectional area of polymer chains¹⁸. The cross-sectional area of polymer chains can be determined from many methods¹⁸⁻²⁰, e.g. directly from the unit-cell parameters or indirectly from empirical correlations. From T_g in *Table 1* and the cohesive energy density E_c in ref. 13, the cross-sectional area of polymer chains for 1,2-polybutadiene, A were determined using¹⁸:

$$T_{\rm g} = 116 + 4.33 E_{\rm c} A \tag{11}$$

where E_c is taken as 72.3 cal cm⁻³ for all 1,2polybutadiene samples¹³. Equation (12) gives a correlation of logarithmic value of monomeric friction coefficient with cross-sectional area of polymer chains, A (in Å²), for 1,2-polybutadiene:

$$\log \zeta_0 = -11.59 + 0.15A \tag{12}$$

with correlation coefficient 0.992. The meaning of (12) is that for the 1,2-polybutadiene series the resistance to driving per monomer unit through the local surroundings at unit velocity depends on the chain flexibility.

The resistance to segment mobility through its local surroundings determines the relaxation process. The relaxation process corresponding to the transition zone of the relaxation spectrum can be described by the shortest relaxation process of Doi-Edwards theory¹⁵. The shortest relaxation time in Doi-Edwards theory is:



Figure 4 Dependence of plateau modulus on the vinyl content for 1,2polybutadienes. Data are taken from ref. 10. Closed circles were measured at 25°C, open circles at 24°C and crosses at 50°C

 Table 2
 Relative values of monomeric friction coefficient for 1,2-polybutadienes

Vinvl content (%)	91	76	53	1/
C_1, C_2 in WLF equation	7.19, 73.11	5.99, 91.25	4.65, 117.8	3.46, 158.3
a (Å)	7.19	7.26	7.36	7.53
$\log[\zeta_0 (g s^{-1})]$	-4.12	- 5.51	-6.28	-8.53
Ne	89	70	52	36
$\log[\tau_{A}(s)]$	-2.89	-4.49	- 5.50	-8.04
d (Å) from equation (3)	36.6	34.1	30.1	23.8
d (Å) from equation (1)	60.8	54.1	47.4	43.5



Figure 5 Dependence of logarithmic shortest relaxation time on the tube diameter in Doi-Edwards theory for 1,2-polybutadienes. Closed circles refer to the tube diameter obtained from equation (3) and open circles from equation (1)

$$\tau_{\rm A} = \frac{1}{6\pi^2} \frac{\zeta_0 a^2}{kT} N_{\rm e}^2 \tag{13}$$

where $N_e = M_e/M_0$ and M_e is the entanglement molecular weight, which is equal to

$$M_{\rm e} = \rho R T / G_{\rm N}^0 \tag{14}$$

The plateau modulus for polybutadiene samples varies with the vinyl content as shown in *Figure 4*. All the data in *Figure 4* are taken from ref. 10; except for the lowest and the highest vinyl contents, all the other values fall reasonably close to the same line. The line in *Figure 4* can be formulated to:

$$G_{\rm N}^0 = 12.42 - 8.65 X_{1,2} \ (10^5 \,{\rm Pa})$$
 (15)

with correlation coefficient -0.986. Combining (15) with (13), the shortest relaxation time of 1,2-polybutadiene samples is obtained and listed in *Table 2*. Those values correspond to τ values at the transition zone in *Figure 3*.

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The shortest relaxation process in Doi-Edwards theory involves equilibration of the chain over the

distance of the tube diameter. This means that the larger the tube diameter, the longer is the relaxation time. The relationship between the shortest relaxation time and the tube diameter for 1,2-polybutadienes is:

$$\log \tau_{\rm A} = -17.24 + 0.38d \tag{16}$$

where the correlation coefficient is 0.992 and the tube diameter is obtained from equation (3). The tube diameter obtained from equation (1) can also be expressed in terms of the shortest relaxation time with a similar formula:

$$\log \tau_{\rm A} = -19.20 + 0.27d \tag{17}$$

with correlation coefficient 0.958. Figure 5 shows τ_A vs. d plots with d obtained from equations (1) and (3).

CONCLUSIONS

The monomeric friction coefficient for 1,2-polybutadiene is related to T_g and the cross-sectional area of polymeric chains. The quantitative relationship between the shortest relaxation time and the tube diameter in Doi– Edwards theory has been found for 1,2-polybutadienes.

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